# High pressure X-ray diffraction study of UC<sub>2</sub> using synchrotron radiation

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#### Abstract

 $UC_2$  has been studied under the general topic of X-ray diffraction analysis of pressure-induced structural transitions in actinide compounds and the determination of their ambient pressure bulk moduli. At 17.6 GPa  $UC_2$  shows a phase transition from a tetragonal to a hexagonal structure. The mean values of the bulk modulus and its pressure derivative are 216 GPa and 2.2 respectively.

# 1. Introduction

Hill [1] has reported that it is possible to associate the magnetic behaviour of actinide compounds with the inter-actinide distance. Below a certain limit of the inter-actinide distance, called the Hill limit, the compounds are non-magnetic which is correlated with a delocalization of the 5f electrons. Above this limit, the compounds are magnetic and the 5f electrons are localized. This Hill criterion has some exceptions but is a good starting point in the study of actinide compounds. Using Hill plots [1], it can be seen that UC<sub>2</sub> falls in the transition region. With increasing pressure, the shortest U–U distance (352 pm at ambient pressure) may decrease sufficiently to induce a structural transition correlated with 5f electron delocalization.

# 2. Experimental details

The uranium dicarbide used was prepared by arc melting from a mixture of uranium metal (purity, 99.7%) and graphite (purity, 99.5%) in the atomic ratio 1:2.5. UC<sub>2</sub> crystallizes in the tetragonal CaC<sub>2</sub>-type structure,

space group *I4/mmm* (no. 139), with two molecules per unit cell. The compound obtained was sealed under a dry atmosphere to avoid reaction with water.

The sample was finely crushed in silicon oil dried over calcium hydride. This same oil was used as the pressure-transmitting medium. The sample was then quickly loaded into an Inconel-gasketed diamond anvil cell of the Syassen-Holzapfel type. A ruby chip was added to the sample for the pressure measurement using the non-linear ruby fluorescence method [2]. The polychromatic X-ray source needed for energydispersive diffraction experiments was synchrotron radiation provided by DESY Hasylab. The Bragg angle was determined from an NaCl diffraction spectrum used as reference. For each pressure step, the  $d_{hkl}$  values deduced from the energy-dispersive spectrum were used to calculate the lattice parameters and the relative volume  $V/V_0$ . The isothermal bulk modulus and its pressure derivative were determined by fitting the  $V/V_0 = f(p)$  data (zero indicates ambient pressure) to several equations of state [3-6]. The lattice parameters determined at ambient pressure in the diamond anvil cell were equal to 551.6(6) pm for a and 600.3(6) pm for c, in agreement with the literature [7].

#### 3. Results and discussion

UC<sub>2</sub> was studied up to 50.1 GPa in 21 steps. Near 17.6 GPa, a stuctural transition occurred as the  $d_{hkl}$  variation shows in Fig. 1. The new phase remains stable under pressure up to the highest pressure reached and, on pressure release, down to approximately 5 GPa. The low pressure structure reappears below this pressure value. To determine the structure of the new phase,

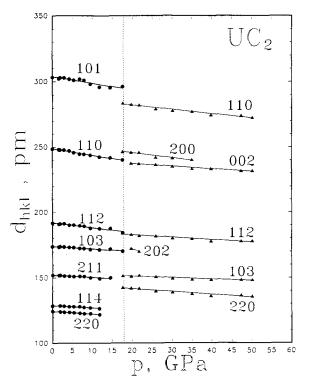


Fig. 1.  $d_{hkl}$  of UC<sub>2</sub> vs. pressure (filled circles for tetragonal low pressure phase and filled triangles for hexagonal high pressure phase).

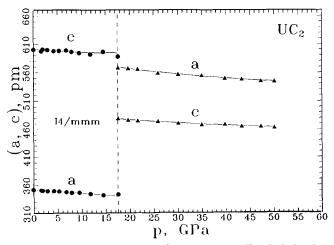


Fig. 2. Lattice parameters of  $UC_2 vs.$  pressure (filled circles for tetragonal low pressure phase and filled triangles for the hexagonal high pressure phase).

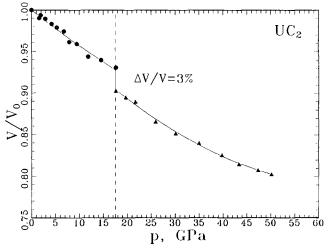


Fig. 3. Relative volume of  $UC_2$  vs. pressure (filled circles for tetragonal low pressure phase and filled triangles for the hexagonal high pressure phase).

TABLE 1. Diffraction data of the hexagonal phase of UC<sub>2</sub> at 19.7 GPa. The reliability factor  $\Sigma |\Delta d| / \Sigma d_{obs} = 0.3\%$  and the lattice parameters are a = 567.4(6) pm and c = 476.7(5) pm

Peak (intensities)	d <sub>obs</sub> (pm)	hkl	d <sub>calc</sub> (pm)	$ \Delta d/d \times 100$
d <sub>1</sub> (w)	282.7	110	283.7	0.35
$d_{2}(s)$	246.1	200	245.7	0.16
$d_{3}(s)$	237.7	002	238.4	0.27
d₄ (m)	182.9	112	182.5	0.20
d <sub>5</sub> (m)	172.3	202	171.1	0.81
$d_6(m)$	151.2	103	151.2	0.02
$d_7$ (w)	142.2	220	141.8	0.27

TABLE 2. Isothermal bulk modulus of tetragonal UC<sub>2</sub> and its pressure derivative determined from several equations of state (Table 3). Average uncertainty is equal to 10 GPa for  $B_0$  and 0.1 for  $B_0'$ 

Equation	B <sub>0</sub> (GPa)	B <sub>0</sub> '	
Murnaghan [4]	216	2.2	
Birch [3]	215	2.4	
Rose et al. [5]	216	2.2	
General MVL [6]	215	2.4	
Holzapfel 0L [6]	216	2.1	
Holzapfel 1L [6]	216	2.0	
Mean	216	2.2	

MVL: Modified Vinet form of Lth order.

H0L: Holzapfel form of Lth order.

H1L: Holzapfel form of Lth order +1 condition compared to H0L.

we made use of two simple mathematical relations observed for the  $d_{hkl}$  values at 19.7 GPa (Table 1):  $d_1/d_7 = 1/2$  and  $d_1/d_2 = (4/3)^{1/2}$ .

From the well-known equations used to calculate the  $d_{hkl}$  values from the lattice parameters and the hkl

TABLE 3. Several equations of state (EOS) used to calculate the bulk modulus  $B_0$  and its pressure derivative  $B_0'$   $(x = (V/V_0)^{1/3} 1 \le k \le L)$ 

Murnaghan [4]	$p = (3/c_2)B_0 x^{-c_2}(1 - x^{-c_2})$ $c_2 = 3B_0' \neq 0$
Birch [3]	$p = 3/2B_0 x^{-7} (1-x^2) (1 + \sum c_{k+1} [(1-x^2)/x^2]^k)$ $c_2 = 3/4 (B_0' - 4)$
Rose et al. [5]	$p = 3B_0(x-1)(\exp(-a^*))[1-0.15a^*+0.05a^{*2}]x^2$ $a^* = (r-r_0)/l \text{ where } r \text{ is Wigner-Seitz radius}$ $l = [\Delta E/(12\pi B_0 r_0)]^{1/2} \text{ where } \Delta E \text{ is cohesive energy}$
General MVL [6]	$p = 3/nB_0 x^{-m}(1-x^n)(\exp[\Sigma c_{k+1}(1-x)^k])$ $p_{FG0} = p_{FG}(Z, V_0), Z = \text{atomic number}$ $c_2 = 3/2[B_0' - (2m-n)/3]$ <i>m</i> and <i>n</i> : adjustable parameters
Holzapfel 0L [6]	$p = 3B_0 x^{-5} (1-x) (\exp[\Sigma c_{k+1} (1-x)^k])$ $c_2 = 3/2 (B_0' - 3)$
Holzapfel 1L [6]	$p = 3B_0 x^{-5} (1-x) [\exp(c_0 (1-x) + x \sum c_{k+1} (1-x)^k)]$ $c_0 = \ln p_{FG0} - \ln B_0$ $c_2 = 3/2 (B_0' - 3) - c_0$

TABLE 4. Linear compressibilities of  $UC_2$  tetragonal lattice parameters [8]

	a	с
β (GPa <sup>-1</sup> )	198(18)×10 <sup>-6</sup>	57.0(40)×10 <sup>-6</sup>
δ (GPa <sup>-2</sup> )	$23.7(13) \times 10^{-6}$	$-25.8(22) \times 10^{-6}$

TABLE 5. Estimated interatomic distances of the  $UC_2$  tetragonal phase at ambient pressure and 17.6 GPa

Atom	Neigh-	Distance (pm)		
	bour	Ambient pressure ( $a = 351.6$ , $c = 600.3$ )	17.6 GPa $(a = 343, c = 588)$	
U2(a)	4U	352	343	
(0,0,0)	8U	390	381	
	2C	233	228	
	8C	258	251	
	2C	367	360	
	8C	422	412	
C4(e)	1C	135	132	
(0,0,2)	4C	299	255	
z = 0.388	4C	352	343	
	4C	376	367	
	8C	390	381	

values, we found that these relations can be simply related to a hexagonal lattice  $(l/d_{hkl}^2 = (4/3)(h^2 + hk + k^2)/a^2 + l^2/c^2$  and l=0). The complete indexing of the pattern observed (Table 1) confirmed this assumption.

From the variation of  $d_{hkl}$  in both phases (Fig. 1), the evolution of the lattice parameters vs. pressure was calculated (Fig. 2). Assuming Z=4 for the hexagonal

unit cell, Fig. 3 shows the pressure dependence of the normalized relative volume for both phases. There is a volume decrease of about 3% at the transition calculated as

$$(V_{\rm hex} - V_{\rm tet})/V_{\rm tet} = -3\%$$

where  $V_{\text{hex}}$  and  $V_{\text{tet}}$  are the unit cell volumes of the hexagonal and tetragonal phases divided by  $Z_{\text{hex}}$  and  $Z_{\text{tet}}$  respectively at the transition pressure.

The isothermal bulk moduli and their pressure derivatives were obtained by fitting several equations of state to the experimental curve of the tetragonal low pressure phase (Tables 2 and 3). Furthermore, assuming hydrostatic conditions, the linear compressibilities of the UC<sub>2</sub> lattice parameters were calculated as proposed by Fietz *et al.* [8] and reported in Table 4.

X-ray diffraction results do not allow us to obtain the atomic positions from the peak intensities, which are not reliable due to possible preferred orientation of the crystals in the polycrystalline sample. Making the approximation that the atomic positions do not change until transition, we can estimate the interatomic distance variation (Table 5). From this table, it seems possible that the first uranium-uranium distance decreases sufficiently to induce a structural change due to the 5f electrons. However, this phase transition may also be due to the decrease in the shortest C-C distance. To obtain more direct information about the electronic behaviour of UC<sub>2</sub> under pressure, spectroscopic measurements, such as optical reflectivity, and resistivity experiments under pressure seem desirable. Furthermore, considering the strong hysteresis observed, it seems that the lattice energies of both structures are similar.

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